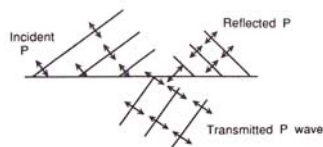
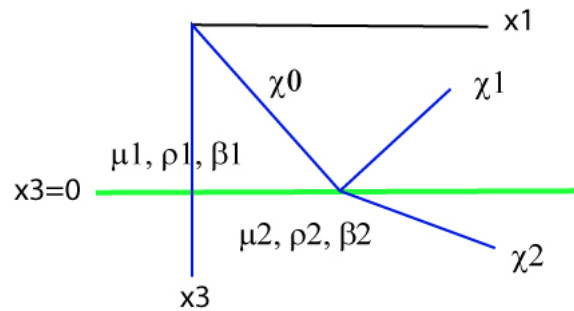
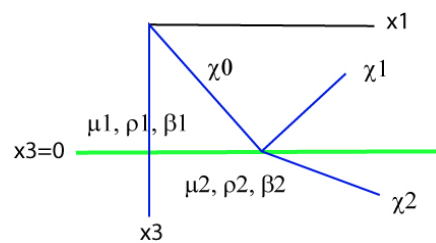


## Wave Solutions in Vertically Varying Media

- Spherical & plane wave solutions considered thus far are valid only within homogeneous space, or smoothly varying media.
- For vertically inhomogeneous media new solutions are found by solving boundary value problems, where there is continuity of traction and displacement across media discontinuities. i.e. stress & displacement are transmitted and surfaces remain in contact.
- We consider reflection and refraction of SH waves at a welded discontinuity.



**FIGURE 3.24** P-wave particle motions for the incident, reflected, and refracted P waves. Note that if this is a solid-solid boundary, the shear stress in the two layers will *not* match at the boundary, requiring the generation of SV motion in both media.



$$\chi_0(\bar{x}, t) = A \cdot e^{i(\omega t - \omega p x_1 - \omega \eta_1 x_{31})}$$

$$\chi_1(\bar{x}, t) = B \cdot e^{i(\omega t - \omega p x_1 + \omega \eta_1 x_{31})}$$

$$\chi_2(\bar{x}, t) = C \cdot e^{i(\omega t - \omega p x_1 - \omega \eta_2 x_{31})}$$

## System of Equations

$$\bar{u} = \nabla\phi + \nabla \times \langle 0, 0, \chi \rangle = -\frac{\partial\chi}{\partial x_1} \hat{x}_2 \quad \text{Helmholtz Equation}$$

$$\sigma_{ij} = \lambda\theta\delta_{ij} + 2\mu\varepsilon_{ij} \quad \sigma_{32} = \mu \frac{\partial u_2}{\partial x_3} = \mu \frac{\partial^2 \chi}{\partial x_1 \partial x_3} \quad \text{Constitutive Equation}$$

$$\left(-\frac{\partial\chi_0}{\partial x_1}\right)_{x_3=0} + \left(-\frac{\partial\chi_1}{\partial x_1}\right)_{x_3=0} = \left(-\frac{\partial\chi_2}{\partial x_1}\right)_{x_3=0} \quad \text{First consider continuity of displacement at } x_3=0 \text{ boundary}$$

$$i\omega p A e^{i(\omega t - \omega p x)} + i\omega p B e^{i(\omega t - \omega p x)} = i\omega p C e^{i(\omega t - \omega p x)}$$

$$A + B = C$$

$$\left(-\mu_1 \frac{\partial^2 \chi_0}{\partial x_1 \partial x_3}\right)_{x_3=0} + \left(-\mu_1 \frac{\partial^2 \chi_1}{\partial x_1 \partial x_3}\right)_{x_3=0} = \left(-\mu_2 \frac{\partial^2 \chi_2}{\partial x_1 \partial x_3}\right)_{x_3=0} \quad \text{Next, consider continuity of traction at } x_3=0$$

$$-(-i\omega p)(-i\omega\eta_1\mu_1)Ae^{i(\omega t - \omega p x_1)} - (-i\omega p)(i\omega\eta_1\mu_1)Be^{i(\omega t - \omega p x_1)} = -(-i\omega p)(-i\omega\eta_2\mu_2)Ce^{i(\omega t - \omega p x_1)}$$

$$A\mu_1\eta_1 - B\mu_1\eta_1 = C\mu_2\eta_2$$

$$A + B = C \quad \text{Equation 1}$$

$$A\mu_1\eta_1 - B\mu_1\eta_1 = C\mu_2\eta_2 \quad \text{Equation 2}$$

First solve:  $(\mu_1\eta_1)(1) + (2)$

$$\begin{aligned} \mu_1\eta_1 A + \mu_1\eta_1 B &= \mu_1\eta_1 C \\ + \\ \mu_1\eta_1 A - \mu_1\eta_1 B &= \mu_2\eta_2 C \\ \hline 2\mu_1\eta_1 A &= (\mu_1\eta_1 + \mu_2\eta_2)C \end{aligned}$$

$$T = \frac{C}{A} = \frac{2\mu_1\eta_1}{(\mu_1\eta_1 + \mu_2\eta_2)}$$

Next solve:  $(-\mu_2\eta_2)(1) + (2)$

$$\begin{aligned} -\mu_2\eta_2 A - \mu_2\eta_2 B &= -\mu_2\eta_2 C \\ + \\ \mu_1\eta_1 A - \mu_1\eta_1 B &= \mu_2\eta_2 C \\ \hline (\mu_1\eta_1 - \mu_2\eta_2)A - (\mu_1\eta_1 + \mu_2\eta_2)B &= 0 \end{aligned}$$

$$R = \frac{B}{A} = \frac{(\mu_1\eta_1 - \mu_2\eta_2)}{(\mu_1\eta_1 + \mu_2\eta_2)}$$

Since  $\eta$  is a function of  $p$  (horizontal slowness),  $\eta_j = \frac{\cos(i)}{\beta_j} = \sqrt{\frac{1}{\beta_j^2} - p^2}$ , the reflection (R), and transmission (T) coefficients depend on the angle of incidence of the wave with respect to the media discontinuity

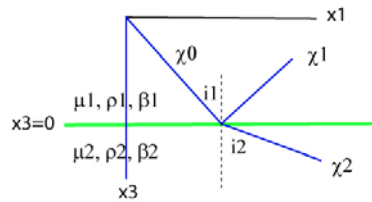
*also remember that  $\mu_1 = \beta_1^2 \rho_1$*

$$T = \frac{2\beta_1\rho_1 \cos(i_1)}{(\beta_1\rho_1 \cos(i_1) + \beta_2\rho_2 \cos(i_2))}$$

$$R = \frac{(\beta_1\rho_1 \cos(i_1) - \beta_2\rho_2 \cos(i_2))}{(\beta_1\rho_1 \cos(i_1) + \beta_2\rho_2 \cos(i_2))}$$

Normal Incidence

$$T = \frac{2\beta_1\rho_1}{(\beta_1\rho_1 + \beta_2\rho_2)} \quad R = \frac{(\beta_1\rho_1 - \beta_2\rho_2)}{(\beta_1\rho_1 + \beta_2\rho_2)}$$



If  $i_1 = i_c$  then  $\eta_2 = 0$  and  $R = 1$  and  $T = 0$

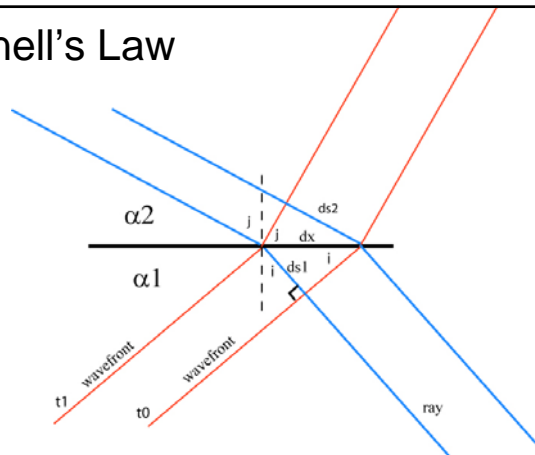
If  $i_1 > i_c$  we have wide-angle (total internal) reflection

$T = 0$

$R$  becomes complex

$$\text{if } i_1 > i_c \text{ then } \left(\frac{1}{\beta_2^2} - p^2\right) < 0 \quad \eta_2 \rightarrow \tilde{\eta}_2 = \pm i \sqrt{p^2 - \frac{1}{\beta_2^2}}$$

## Snell's Law



$$dt = \frac{ds_1}{\alpha_1} = \frac{dx \cdot \sin(i)}{\alpha_1} = \frac{dx \cdot \sin(j)}{\alpha_2}$$

$$\frac{dt}{dx} = \frac{\sin(i)}{\alpha_1} = \frac{\sin(j)}{\alpha_2} = p$$

## Snell's Law is General

$$\frac{dt}{dx} = p = \frac{\sin(i)}{\alpha_1} = \frac{\sin(j)}{\alpha_2} = \frac{\sin(k)}{\beta_2} = \text{etc.}$$

$$\text{If } \alpha_2 > \alpha_1 \quad j > i$$

$$\text{If } \alpha_2 < \alpha_1 \quad j < i$$

Behavior of  $R_{\text{SH}}(p)$  for  $p \geq 1/\beta_2$ ,  $i > i_c$

$$R(p) \Big|_{p \geq 1/\beta_2} = \frac{\mu_1 n_1 - i \mu_2 \tilde{n}_2}{\mu_1 n_1 + i \mu_2 \tilde{n}_2}, \quad \tilde{n}_2 = [p^2 - 1/\beta_2^2]^{1/2}$$

$R(p)$  can  $\therefore$  be written in terms of its amplitude and phase, where

$$\|R(p)\| = \frac{[\mu_1^2 n_1^2 + \mu_2^2 \tilde{n}_2^2]^{1/2}}{[\mu_1^2 n_1^2 + \mu_2^2 \tilde{n}_2^2]^{1/2}} = 1$$

$$\text{Phase } \phi = 2 \cdot \tan^{-1} \left( \frac{\mu_2 \tilde{n}_2}{\mu_1 n_1} \right)$$

$$\therefore R_{cp} \Big|_{p > 1/\beta_2} = e^{i\phi}$$

$$\text{and } x_{\uparrow} = 1 \cdot c \quad i\omega(t - px_1 - n_1 x_3 + \frac{\phi}{\omega})$$

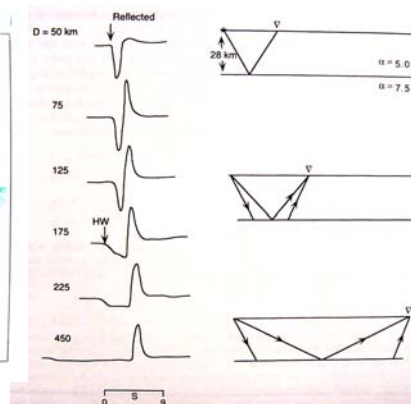
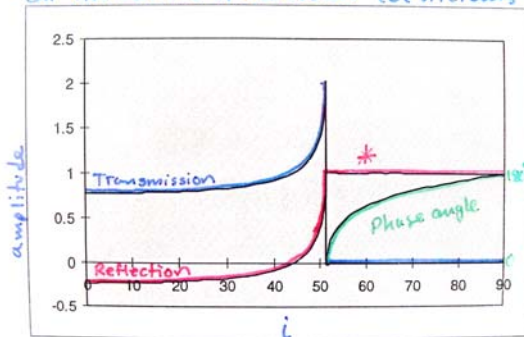
Thus for a post-critically reflected wave  $\|R_{cp}\| = 1$  indicating that the wave is totally reflected however a phase shift is introduced which produces earlier arrivals since

$$t - px_1 - n_1 x_3 + \frac{\phi}{\omega} = c \quad t = px_1 + n_1 x_3 - \frac{\phi}{\omega}$$

assuming  $c=0$

## SH R&T PLOT

SH Reflection & Transmission Coefficients



$$T_{SH}(p) \Big|_{p \gg 1/\beta_2} = \left( \frac{2\mu_1\mu_2}{\mu_1\mu_2 + i\mu_2\tilde{\mu}_2} \right) \cdot \left( \frac{\mu_1\mu_2 - i\mu_2\tilde{\mu}_2}{\mu_1\mu_2 - i\mu_2\tilde{\mu}_2} \right)$$

$$= \frac{2\mu_1^2\mu_2^2}{\mu_1^2\mu_2^2 + \mu_2^2\tilde{\mu}_2^2} - \frac{2i\mu_1\mu_2\mu_2\tilde{\mu}_2}{\mu_1^2\mu_2^2 + \mu_2^2\tilde{\mu}_2^2}$$

For <sup>^</sup> transmitted waves  
 $p \gg 1/\beta_2$   $p$  exists only at  $p = 1/\beta_2$

thus

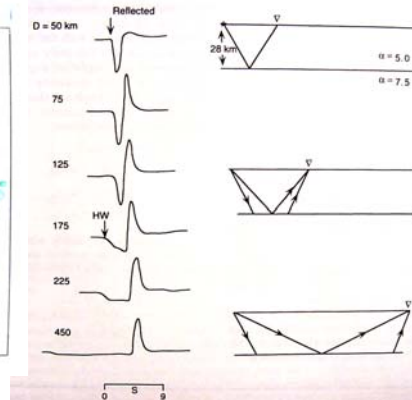
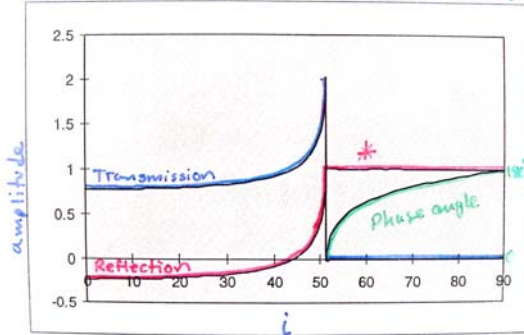
$$\chi_2(\bar{x}, t) = T \left( \frac{1}{\beta_2} \right) e^{i\omega(t - px_1 \pm i\tilde{\mu}_2 x_3)}$$

$$= T(p) e^{-\omega\tilde{\mu}_2|x_3|} e^{i\omega(t - px_1)}$$

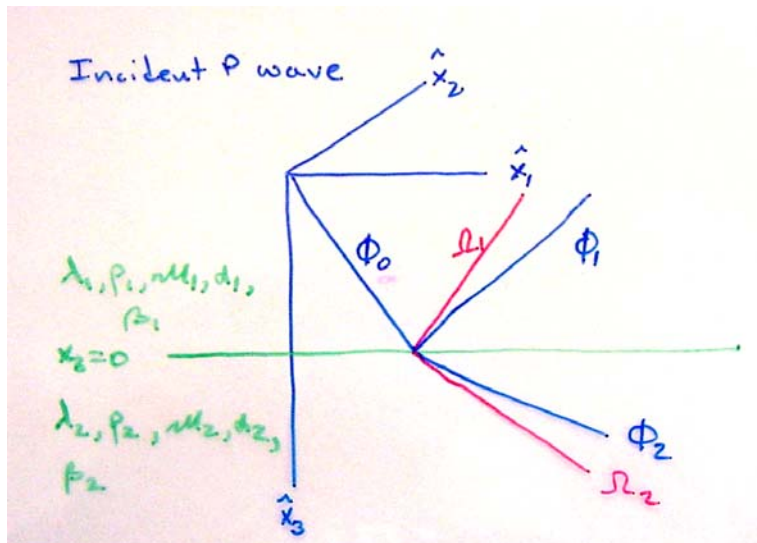
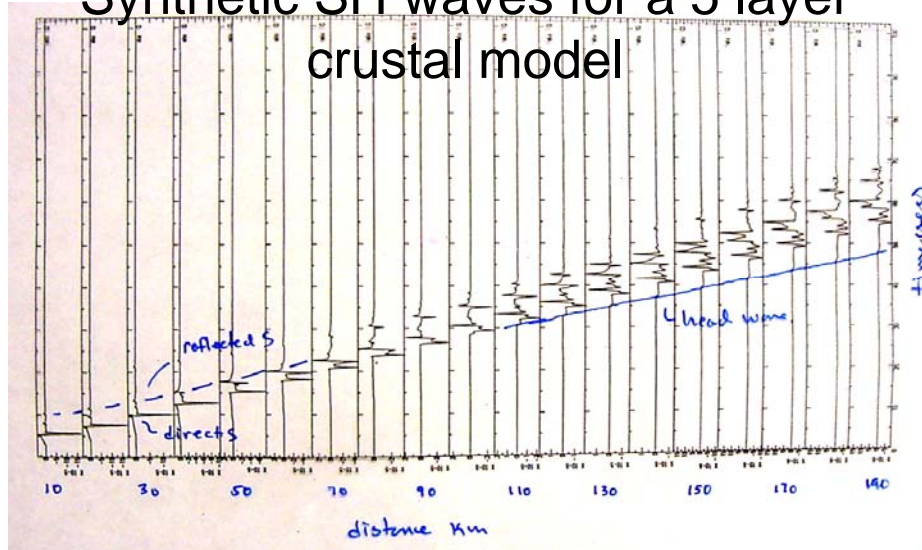
$$= 2e^{i\omega(t - x_1/\beta_2)} \quad x_3 = 0$$

## SH R&T PLOT

SH Reflection & Transmission Coefficients



# Synthetic SH waves for a 5 layer crustal model





$$\phi_0 = A e^{i\omega(t - \rho x_1 - \eta_{\alpha_1} x_3)}$$

$$\phi_1 = B e^{i\omega(t - \rho x_1 + \eta_{\beta_1} x_3)} \quad \Omega_1 = D e^{i\omega(t - \rho x_1 + \eta_{\beta_1} x_3)}$$

$$\phi_2 = C e^{i\omega(t - \rho x_1 - \eta_{\beta_2} x_3)} \quad \Omega_2 = E e^{i\omega(t - \rho x_1 - \eta_{\beta_2} x_3)}$$

$$\bar{\sigma}_{31} = \mu \left( \frac{\partial u_3}{\partial x_1} + \frac{\partial u_1}{\partial x_3} \right)$$

$$\bar{\sigma}_{33} = \lambda \Theta + 2\mu \frac{\partial u_3}{\partial x_3}$$

$$\bar{u} = \left\langle \frac{\partial \phi}{\partial x_1}, \frac{\partial \phi}{\partial x_2}, \frac{\partial \phi}{\partial x_3} \right\rangle + \nabla \times \langle 0, \Omega, 0 \rangle$$

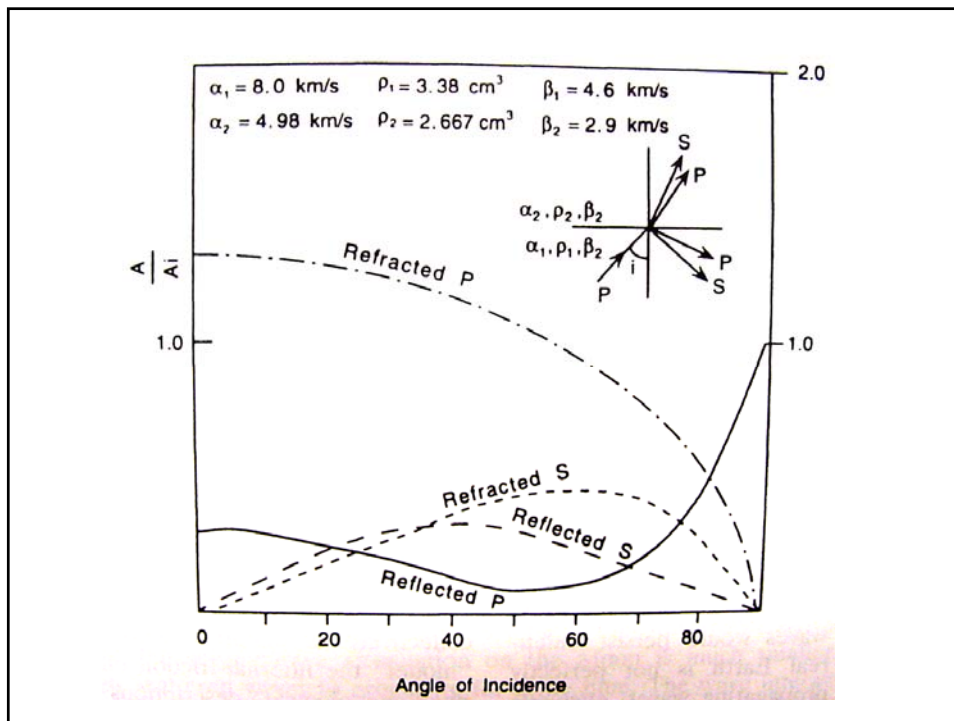
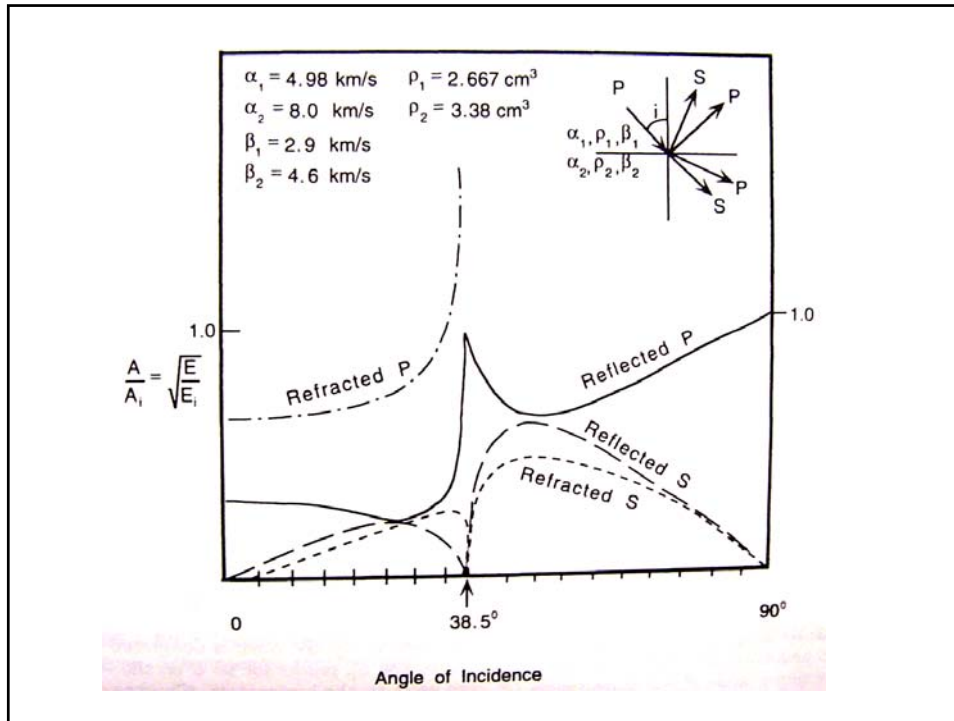
$$= \left( \frac{\partial \phi}{\partial x_1} - \frac{\partial \Omega}{\partial x_3} \right) \hat{x}_1 + \left( \frac{\partial \phi}{\partial x_3} + \frac{\partial \Omega}{\partial x_1} \right) \hat{x}_3$$

Solve boundary value problem for continuity of stress across  $\hat{x}_3$  surface ( $\bar{\sigma}_{31}, \bar{\sigma}_{33}$ ) and displacement in  $\hat{x}_1, \hat{x}_3$  plane.

**TABLE 3.1** Displacement, Reflection and Transmission Coefficients

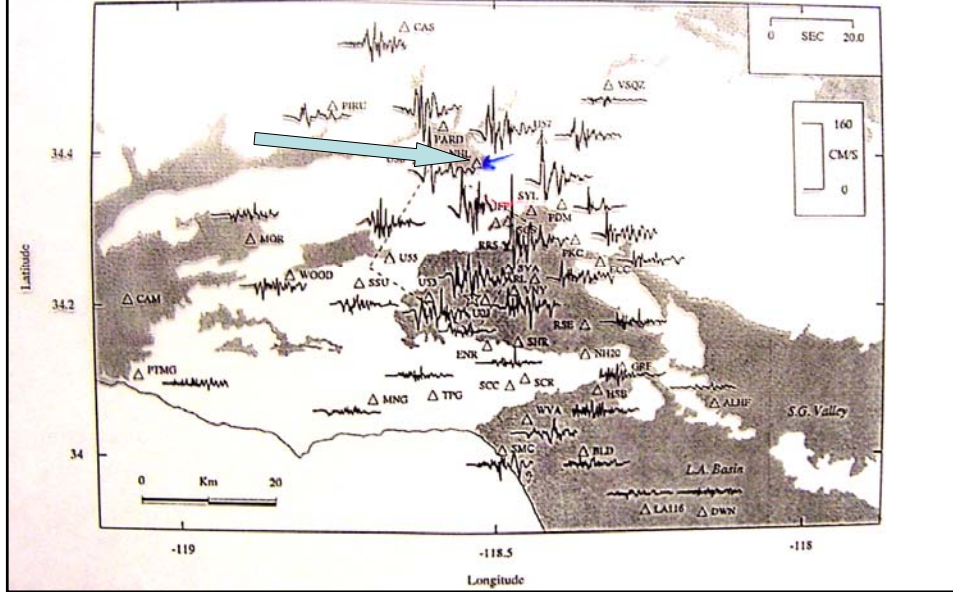
Coefficient	Formula
Solid-free surface (P-SV)	
$R_{PP}$	$[-(1/\beta^2) - 2\rho^2]^2 + 4\rho^2\eta_\alpha\eta_\beta / A$
$R_{PS}$	$(4(\alpha/\beta)\rho\eta_\alpha[(1/\beta^2) - 2\rho^2]) / A$
$R_{SP}$	$(4(\beta/\alpha)\rho\eta_\beta[(1/\beta^2) - 2\rho^2]) / A$
$R_{SS}$	$[-(1/\beta^2) - 2\rho^2]^2 + 4\rho^2\eta_\alpha\eta_\beta / A$
$R_{SS}(SH)$	1
Solid-solid (P-SV)	
$R_{PP}$	$[(b\eta_{\alpha_1} - c\eta_{\alpha_2})F - (a + d\eta_{\alpha_1}\eta_{\beta_2})Hp^2] / D$
$R_{PS}$	$-[2\eta_{\alpha_1}(ab + cd\eta_{\alpha_2}\eta_{\beta_2})\rho(\alpha_1/\beta_1)] / D$
$T_{PP}$	$[2\rho_1\eta_{\alpha_1}F(\alpha_1/\alpha_2)] / D$
$T_{PS}$	$[2\rho_1\eta_{\alpha_1}Hp(\alpha_1/\beta_2)] / D$
$R_{SS}$	$-(b\eta_{\beta_1} - c\eta_{\beta_2})E - (a + b\eta_{\alpha_2}\eta_{\beta_1})Gp^2 / D$
$R_{SP}$	$-[2\eta_{\beta_1}(ab + cd\eta_{\alpha_2}\eta_{\beta_2})\rho(\beta_1/\alpha_1)] / D$
$R_{SS}(SH)$	$\frac{\mu_1\eta_{\beta_1} - \mu_2\eta_{\beta_2}}{\mu_1\eta_{\beta_1} + \mu_2\eta_{\beta_2}}$
$T_{SS}(SH)$	$\frac{2\mu_1\eta_{\beta_1}}{\mu_1\eta_{\beta_1} + \mu_2\eta_{\beta_2}}$
$a = \rho_2(1 - 2\beta_2^2\rho^2) - \rho_1(1 - 2\beta_1^2\rho^2)$	$E = b\eta_{\alpha_1} + c\eta_{\alpha_2}$
$b = \rho_2(1 - 2\beta_2^2\rho^2) - 2\rho_1\beta_1^2\rho^2$	$F = b\eta_{\beta_1} + c\eta_{\beta_2}$
$c = \rho_1(1 - 2\beta_1^2\rho^2) + 2\rho_2\beta_2^2\rho^2$	$G = a - d\eta_{\alpha_1}\eta_{\beta_2}$
$d = 2(\rho_2\beta_2^2 - \rho_1\beta_1^2)$	$H = a - d\eta_{\alpha_2}\eta_{\beta_1}$
$D = EF + GHp^2$	
$A = [(1/\beta^2) - 2\rho^2]^2 + 4\rho^2\eta_\alpha\eta_\beta$	

Downward incident  
What is the form for upward Incident?



# Example – Northridge Earthquake

Recorded Ground Velocities



D. J. Wald, T. H. Heaton, and K. W. Hudnut

Table I  
Northridge Regional Velocity Structure

A. Rock Stations				
$V_p$ (km/sec)	$V_s$ (km/sec)	Density (g/cm)	Thickness (km)	Depth (km)
1.9	1.0	2.1	0.5	0.0
4.0	2.0	2.4	1.0	0.5
5.5	3.2	2.7	2.5	1.5
6.3	3.6	2.8	23.0	4.0
6.8	3.9	2.9	13.0	27.0
7.8	4.5	3.3		40.0

B. Soil Stations				
$V_p$ (km/sec)	$V_s$ (km/sec)	Density (g/cm)	Thickness (km)	Depth (km)
1	0.8	1.7	0.1	0.0
2	1.2	1.8	0.2	0.1
3	1.9	2.1	0.2	0.3
4	4.0	2.4	1.0	0.5
5	5.5	2.7	2.5	1.5
6	6.3	2.8	23.0	4.0
7	6.8	2.9	13.0	27.0
8	7.8	3.3		40.0

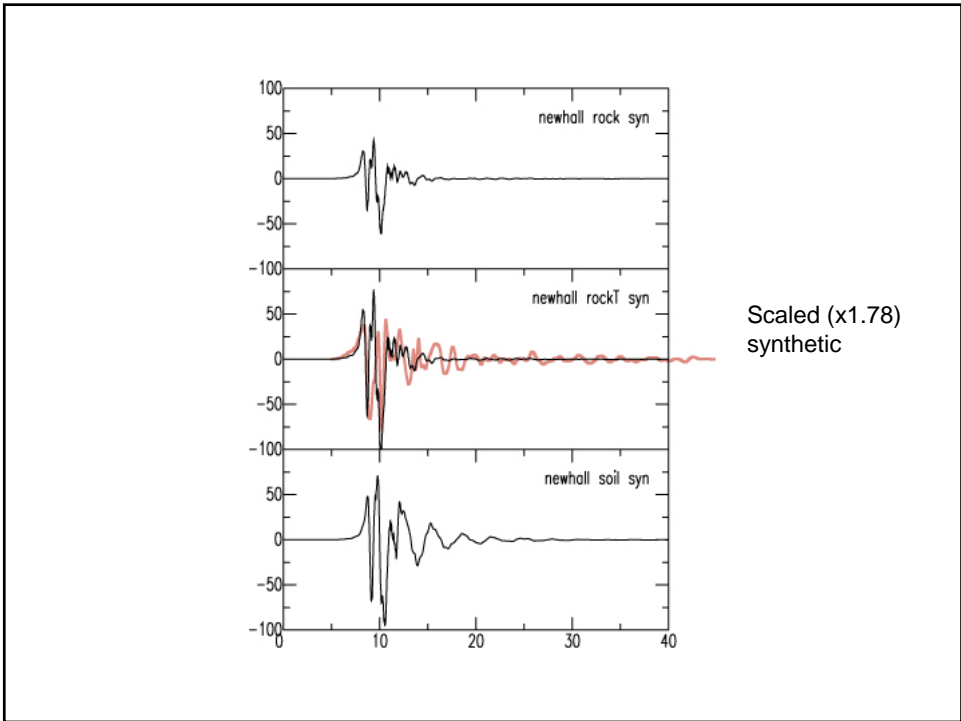
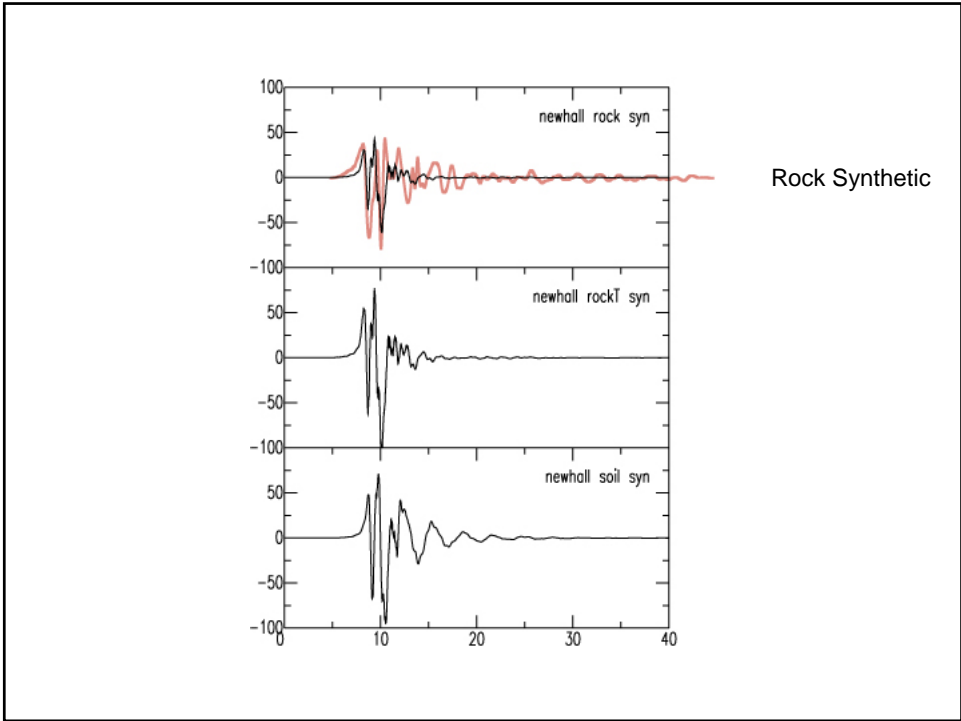
Rock

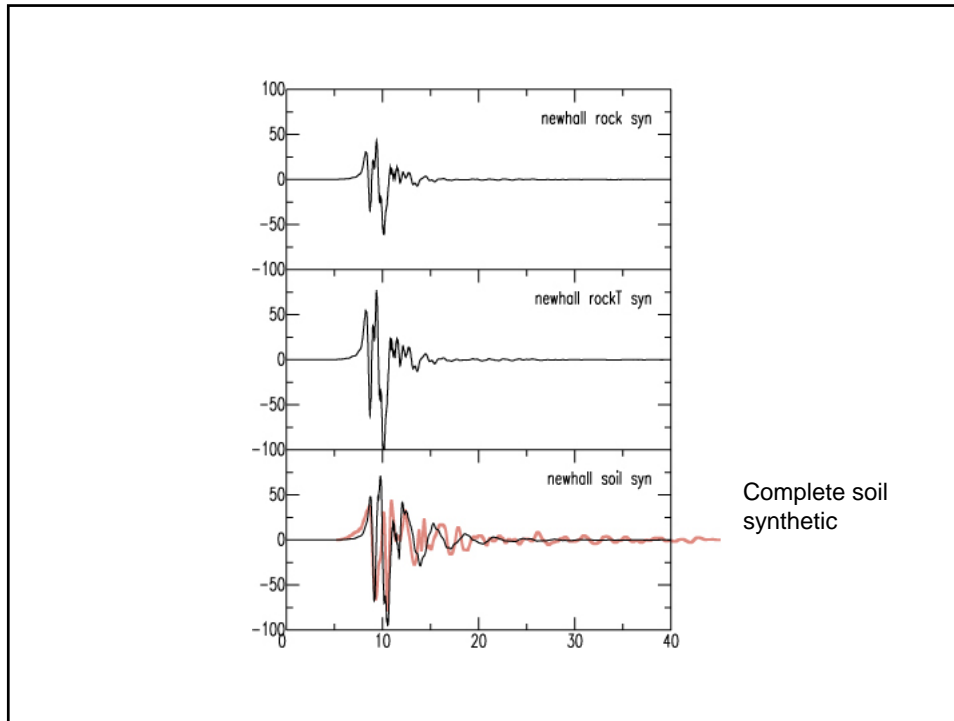
$T_{21}$   
 $T_{32}$

$$T_{32} = 1.40$$

$$T_{21} = 1.27$$

$$T = T_{32} \cdot T_{21} = 1.78$$





$R_{SH}$  @ free surface

$\sigma_{32}$

from the continuity of stress equation

$\mu_1 \frac{\partial^2 \chi_2}{\partial x_1 \partial x_3} \Big|_{x_3=0} + \mu_2 \frac{\partial^2 \chi_2}{\partial x_1 \partial x_3} \Big|_{x_3=0} = 0$

$\mu_1 (-i\omega\rho)(-i\omega m_1)A + \mu_2 (-i\omega\rho)(i\omega m_1) = 0$

$-A + B = 0$

$A = B$

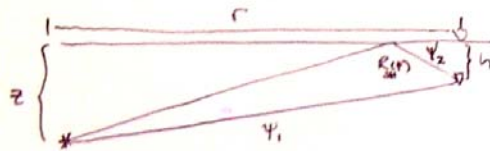
$R_{SH} = 1$  @ free surface

$\sigma_{32} = \phi$

$\sigma_{32} = \frac{\partial^2 \chi_2}{\partial x_1 \partial x_3} \mu$

### Free Surface Focussing

- Consider the following geometry for an incident SH plane wave



- The total displacement potential is then

$$\begin{aligned} \Psi(\vec{x}, t) &= \Psi_1(\vec{x}, t) + \Psi_2(\vec{x}, t) \\ &= A e^{i\omega(t - px_1 + \pi(z-h))} + A R_{SH} e^{i\omega(t - px_1 + \pi z - \pi h)} \\ \lim_{h \rightarrow 0} \Psi(\vec{x}, t) &= (A + A R_{SH}) e^{i\omega(t - px_1 + \pi z)} \\ &= \underline{\underline{2A}} e^{i\omega(t - px_1 + \pi z)} \end{aligned}$$

### Example: Richmond Field Station

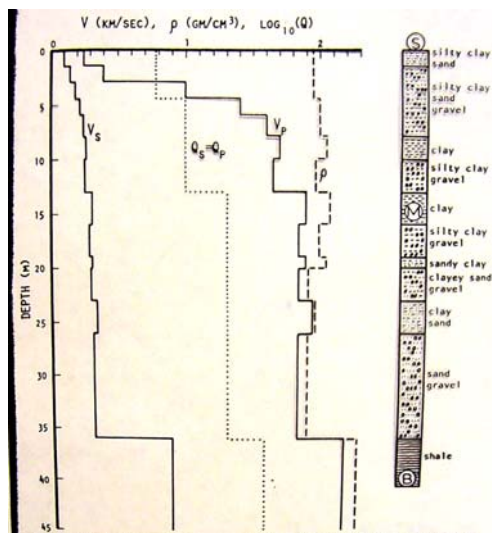


FIG. 2. The depth profiles of velocities, quality factors, and density which were estimated for the vicinity of the vertical seismic array. The section on the right was derived from drill logs and drill cores. The location of seismic sensors at the bottom of the hole, mid-hole, and surface are denoted by B, M, and S, respectively.

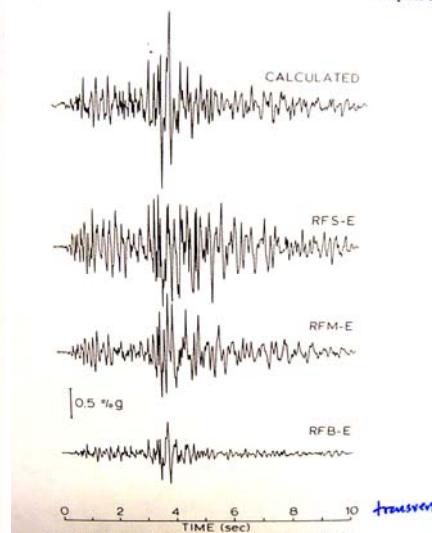
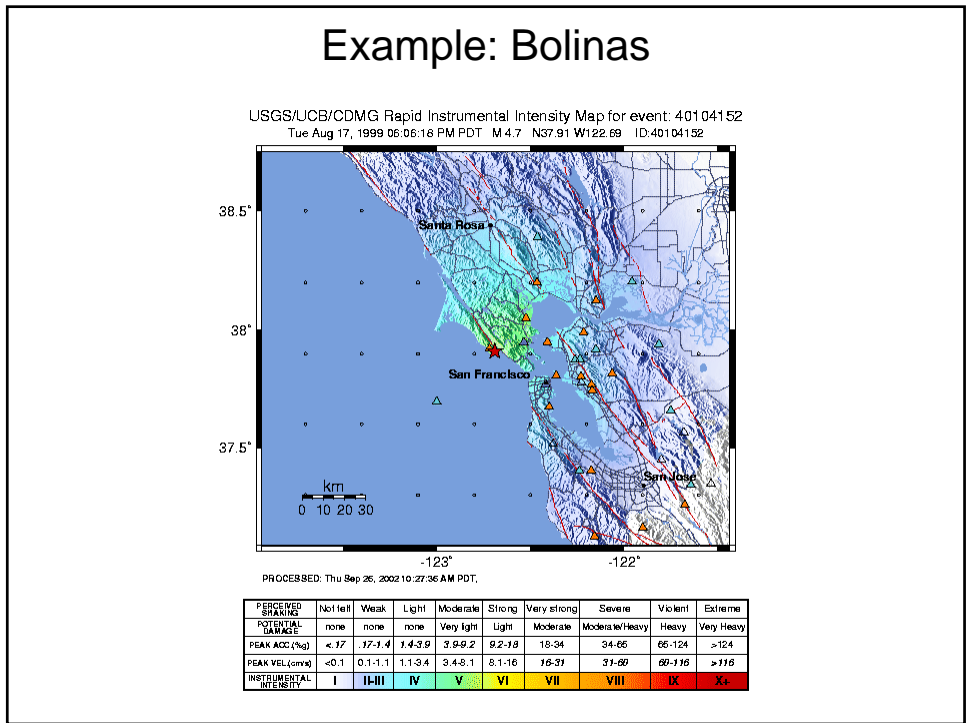
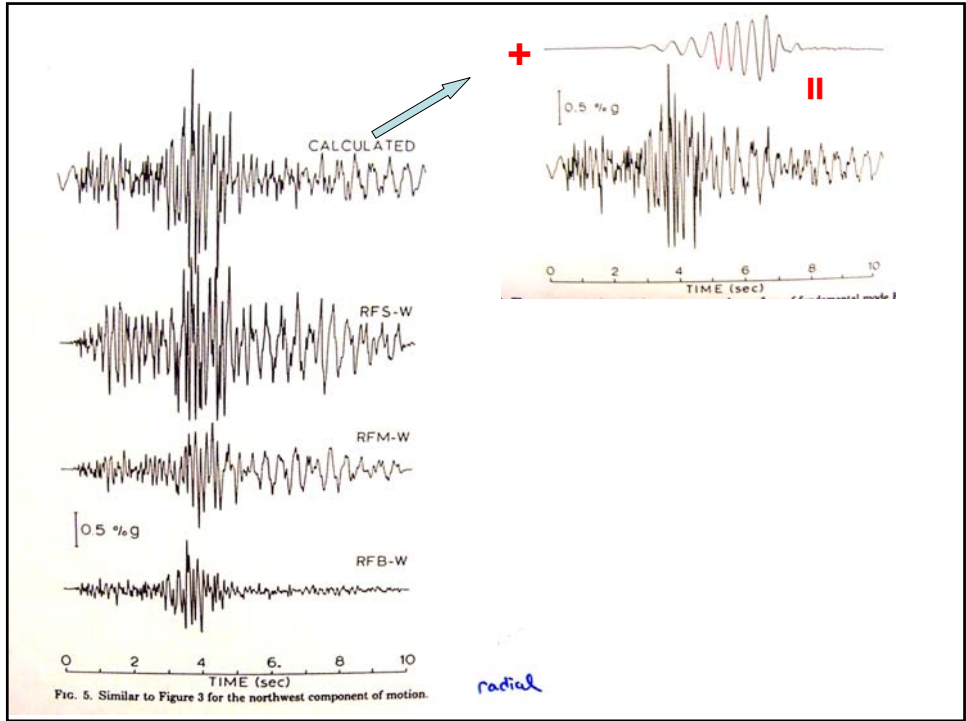
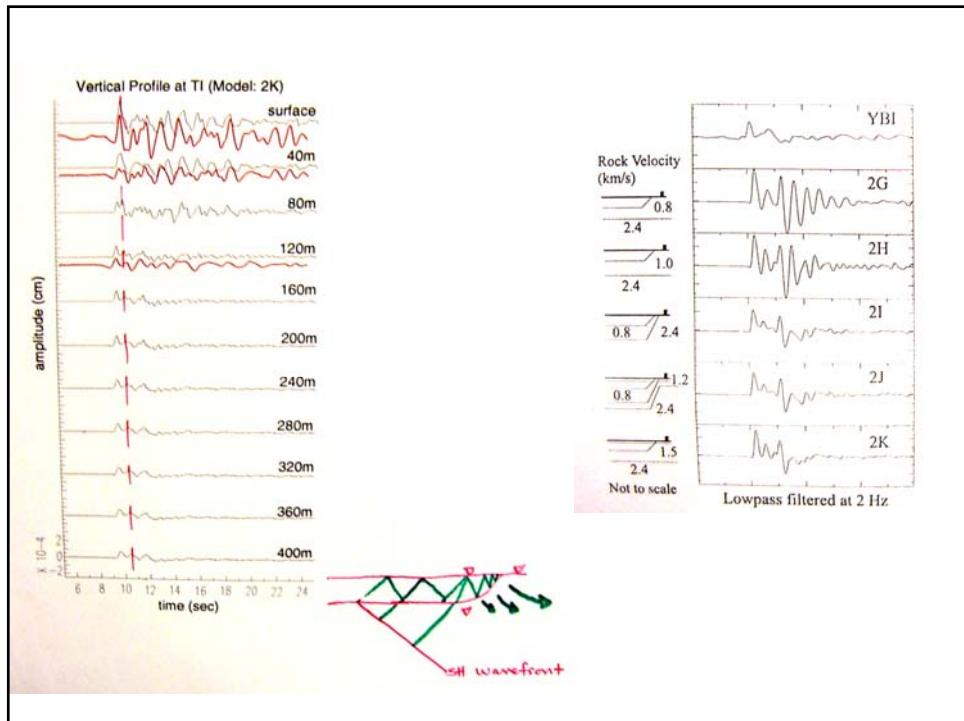
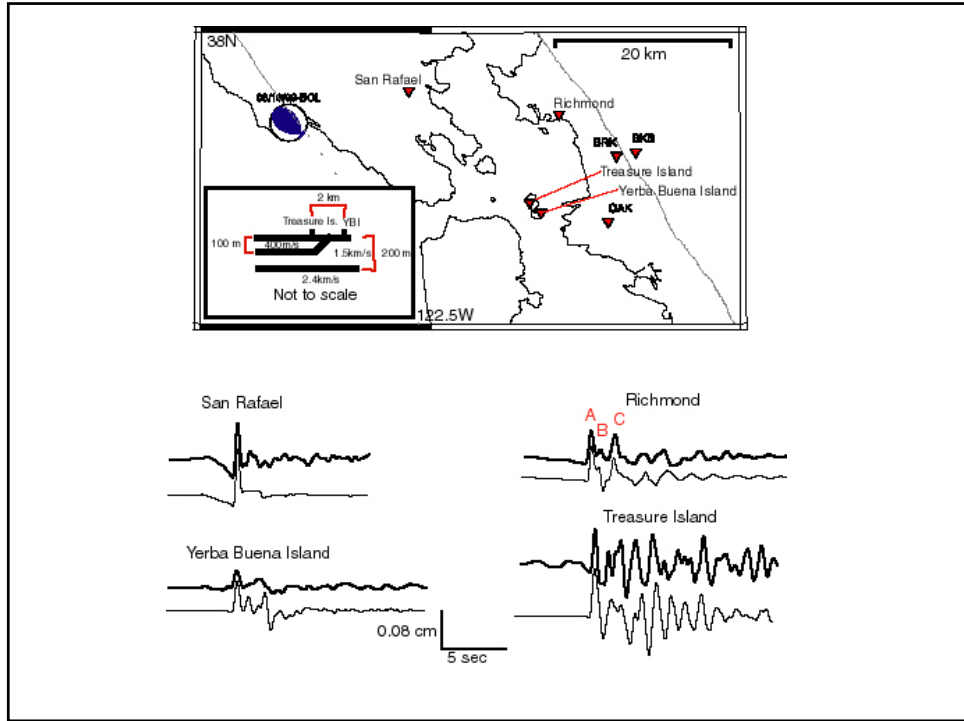


FIG. 4. Similar to Figure 3 for the northeast component of motion.









Energy / Energy Flux should be conserved

$R^2 + T^2 \neq 1$  giving the appearance that energy is not conserved

The fact that energy is carried to and away from an interface at different angles needs to be considered

$$E = \left[ \frac{1}{2} \rho \left( \frac{du}{dt} \right)^2 + \frac{1}{2} \mu \left( \frac{du}{dx} \right)^2 \right] \delta A \delta x$$

$$E = \left[ \frac{1}{2} \rho \left( \frac{du}{dt} \right)^2 + \frac{1}{2} \mu \left( \frac{du}{dx} \right)^2 \right] \delta A \delta x$$

assuming for SH that  $u = B \cos(\omega t - kx)$   
where  $\frac{\omega}{k} = v$

$$\frac{du}{dt} = -B\omega \sin(\omega t - kx) \quad \& \quad \frac{du}{dx} = Bk \sin(\omega t - kx)$$

$$\therefore \frac{du}{dx} = -\frac{1}{v} \frac{du}{dt}$$

then

$$E = \rho \left( \frac{du}{dt} \right)^2 \delta A \delta x = \rho v \left( \frac{du}{dt} \right)^2 \delta A \delta t$$

the energy flux is then

$$I = \rho v \left( \frac{du}{dt} \right)^2 \quad \frac{\text{kg}}{\text{m}^3} \cdot \frac{\text{m}}{\text{s}} \cdot \frac{\text{m}^2}{\text{s}^2} \rightarrow \frac{(\text{N}\cdot\text{m})}{\text{m}^2\cdot\text{s}}$$

Consider the interface at  $x_2 = 0$

$$E = \rho_1 \beta_1 \left( \frac{\partial u}{\partial t} \right)^2 (\delta S \cos i) \delta t$$

$$= \rho_1 \beta_1 \frac{\beta \cos i}{\beta_1} \beta^2 \omega^2 \sin^2(\omega t - kx) \delta t \delta S$$

$$= \mu_1 \eta_1 \beta^2 \omega^2 \sin^2(\omega t - kx) \delta t \delta S$$

$$E_{inc.} = E_{refl.} + E_{trans.}$$

$$\mu_1 \eta_1 = \mu_1 \eta_1 R^2 + \mu_2 \eta_2 T^2$$

$$1 = R^2 + \frac{\mu_2 \eta_2}{\mu_1 \eta_1} T^2$$

### Reflection & Transmitted Energy - Waves Incident from Above

Energy Partitioning in SH Waves

$\rho_1 = 2000$     $\rho_2 = 2670$     $\beta_1 = 2.0$     $\beta_2 = 3.5$    SI Units

$\mu_1 = \beta_1^2 \rho_1$     $\mu_2 = \beta_2^2 \rho_2$

$\eta_1 = 8 \cdot 10^3$     $\eta_2 = 3.271 \cdot 10^4$

$\eta_1(p) = \frac{1}{\sqrt{\beta_1^2 - p^2}}$     $\eta_2(p) = \frac{1}{\sqrt{\beta_2^2 - p^2}}$

$$R(p) = \frac{\mu_1 \eta_1(p) - \mu_2 \eta_2(p)}{\mu_1 \eta_1(p) + \mu_2 \eta_2(p)} \quad T(p) = \begin{cases} \frac{2 \mu_1 \eta_1(p)}{\mu_1 \eta_1(p) + \mu_2 \eta_2(p)} & \text{if } p \leq \frac{1}{\beta_2} \\ 0.0 & \text{otherwise} \end{cases}$$

Transmission Energy Coefficient  $E(p) = \frac{\mu_2 \eta_2(p)}{\mu_1 \eta_1(p)}$  i.e. Ray area normalization factor

$N = 899$   
 $i := 0..N$   
 $\theta_i := \frac{\pi}{180} i \cdot 0.1$   
 $p_i := \frac{\sin(\theta_i)}{\beta_1}$

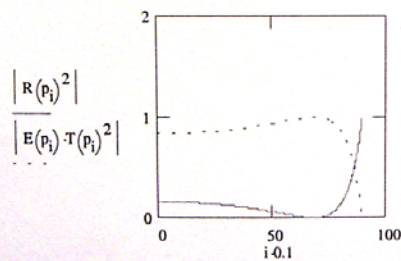
$$\Phi_i = \begin{cases} 0 & \text{if } p_i < \frac{1}{\beta_2} \\ \arg(R(p_i)) - 1 & \text{otherwise} \end{cases}$$

## Reflection & Transmitted Energy - Waves Incident from Below

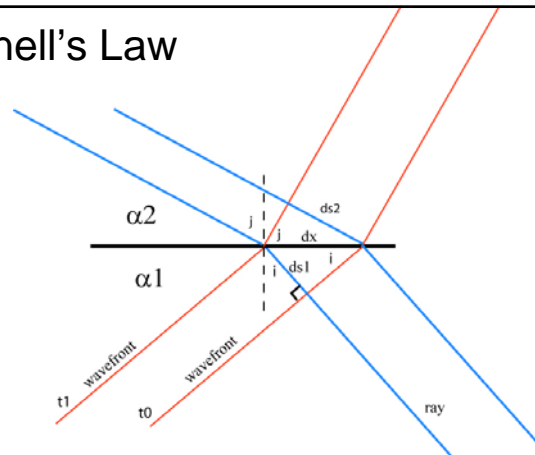
$$R(p) := \frac{(\mu_2 \cdot \eta_2(p)) - \mu_1 \cdot \eta_1(p)}{\mu_1 \cdot \eta_1(p) + \mu_2 \cdot \eta_2(p)} \quad T(p) := \frac{2 \cdot \mu_2 \cdot \eta_2(p)}{\mu_1 \cdot \eta_1(p) + \mu_2 \cdot \eta_2(p)}$$

$$E(p) := \frac{\mu_1 \cdot \eta_1(p)}{\mu_2 \cdot \eta_2(p)}$$

$$p_i := \frac{\sin(\theta_i)}{\beta_2}$$



## Snell's Law



$$dt = \frac{ds_1}{\alpha_1} = \frac{dx \cdot \sin(i)}{\alpha_1} = \frac{dx \cdot \sin(j)}{\alpha_2}$$

$$\frac{dt}{dx} = \frac{\sin(i)}{\alpha_1} = \frac{\sin(j)}{\alpha_2} = p$$

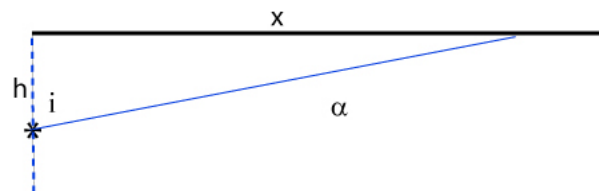
### Snell's Law is General

$$\frac{dt}{dx} = p = \frac{\sin(i)}{\alpha_1} = \frac{\sin(j)}{\alpha_2} = \frac{\sin(k)}{\beta_2} = \text{etc.}$$

$$\text{If } \alpha_2 > \alpha_1 \quad j > i$$

$$\text{If } \alpha_2 < \alpha_1 \quad j < i$$

### Direct Wave



$$\tau = px + \eta h$$

$$= \frac{x}{R\alpha} x + \frac{h}{R\alpha} h$$

$$= \frac{(x^2 + h^2)}{R\alpha} = \frac{\sqrt{x^2 + h^2}}{\alpha} = \frac{R}{\alpha}$$

## Alternative Solution Technique

Suppose  $x$ ,  $h$ , and  $\alpha$  are given. How can we solve for  $p$ , arrival time ( $T$ ) and the reflection transmission coefficients?

Shooting rays

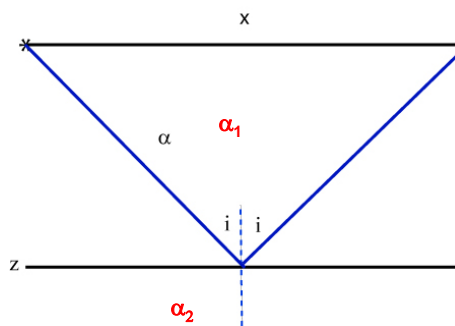
$$x_{\text{given}} = h \cdot \tan(i) \quad \text{Find value of } i \text{ that gives desired distance}$$

$$p = \frac{\sin(i)}{\alpha} \quad \text{Calculate the ray parameter}$$

$$\tau = px + \eta h \quad \text{Calculate arrival time and other parameters (e.g. transmission & reflection coefficients)}$$

This seems more difficult, but there is a clear advantage in more complicated problems.

### Reflection



$$\tau = px + 2\eta z$$

$$= \frac{x}{2R\alpha} x + 2 \frac{z}{R\alpha} z$$

$$= \frac{x^2 + 4z^2}{2\alpha \sqrt{\frac{x^2}{4} + z^2}} = \frac{\sqrt{x^2 + 4z^2}}{\alpha}$$

$$\tau = \frac{2R}{\alpha} = \frac{2\sqrt{\frac{x^2}{4} + z^2}}{\alpha}$$

$$= \frac{\sqrt{x^2 + 4z^2}}{\alpha}$$

## Alternative Method

$$x_{\text{given}} = 2h \cdot \tan(i) \quad \text{Find value of } i \text{ that gives desired distance}$$

$$p = \frac{\sin(i)}{\alpha} \quad \text{Calculate the ray parameter}$$

$$\tau = px + 2\eta h \quad \text{Calculate arrival time and other parameters (e.g. transmission & reflection coefficients)}$$

## Head wave

Arrival Time

$$T = px + 2\eta z$$

$$= \frac{x}{\alpha_2} + 2 \left[ \frac{1}{\alpha_1^2} - \frac{1}{\alpha_2^2} \right]^{1/2} z$$

Critical Distance

$$\tan(i) = \frac{p}{\eta} = \frac{x}{2z}$$

$$x_c = 2z \frac{p}{\eta} = 2z \frac{1}{\beta_2} \frac{1}{\left[ \frac{1}{\beta_1^2} - \frac{1}{\beta_2^2} \right]^{1/2}}$$

$$= 2z \frac{\beta_1}{\sqrt{\beta_2^2 - \beta_1^2}}$$

Cross-over distance

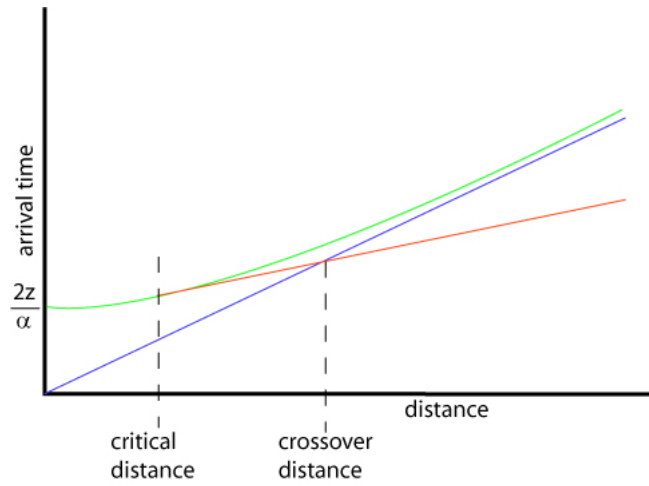
$$\frac{x}{\alpha_1} = \frac{x}{\alpha_2} + 2 \left[ \frac{1}{\alpha_1^2} - \frac{1}{\alpha_2^2} \right]^{1/2} z$$

$$\frac{x^2 (\alpha_2 - \alpha_1)^2}{\alpha_1^2 \alpha_2^2} = 4z^2 \frac{(\alpha_2^2 - \alpha_1^2)}{\alpha_1^2 \alpha_2^2}$$

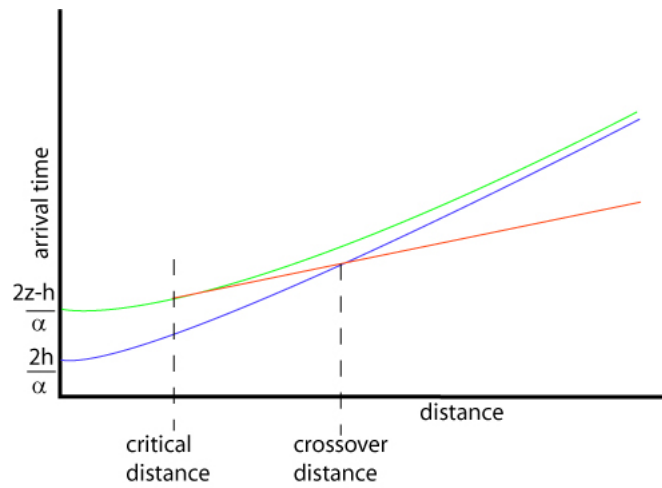
$$x^2 (\alpha_2 - \alpha_1)^2 = 4z^2 (\alpha_2 + \alpha_1)(\alpha_2 - \alpha_1)$$

$$x_c = 2z \sqrt{\frac{(\alpha_2 + \alpha_1)}{(\alpha_2 - \alpha_1)}}$$

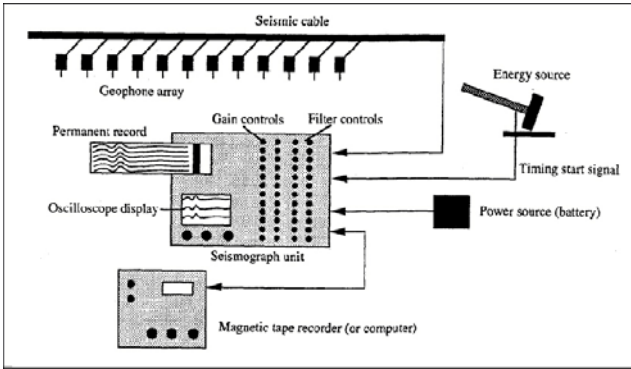
## Typical Traveltime Curve



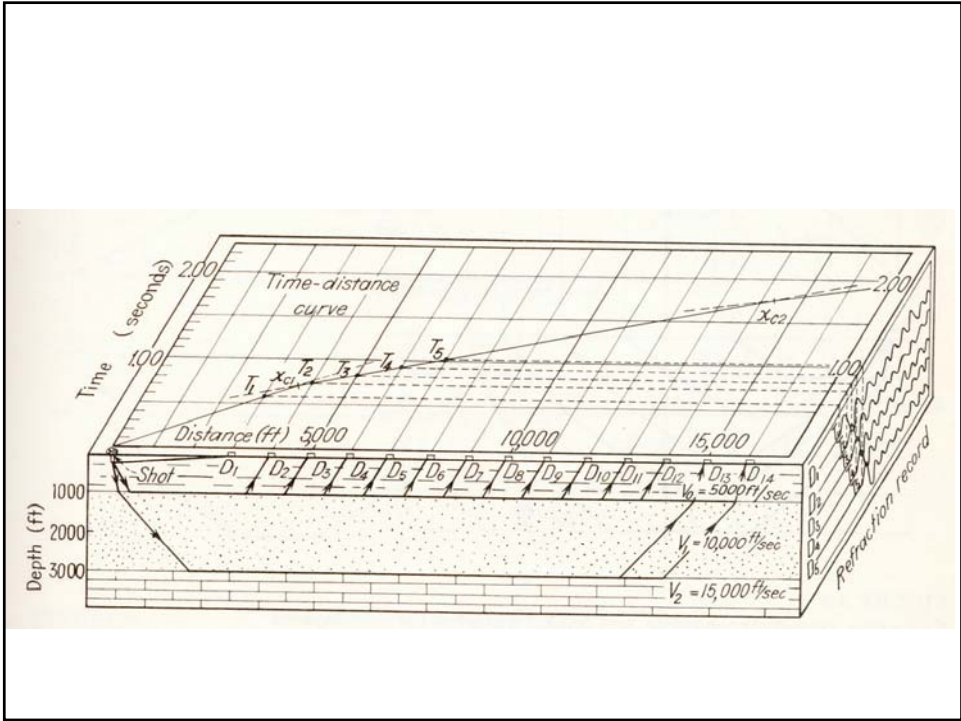
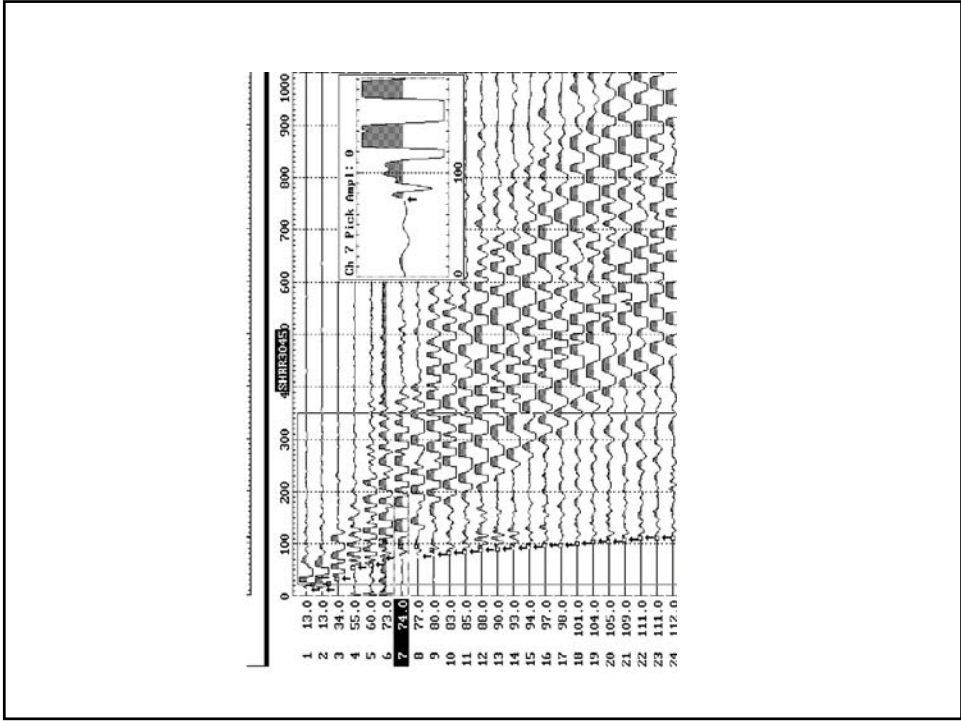
How does it change for a source not at the surface?

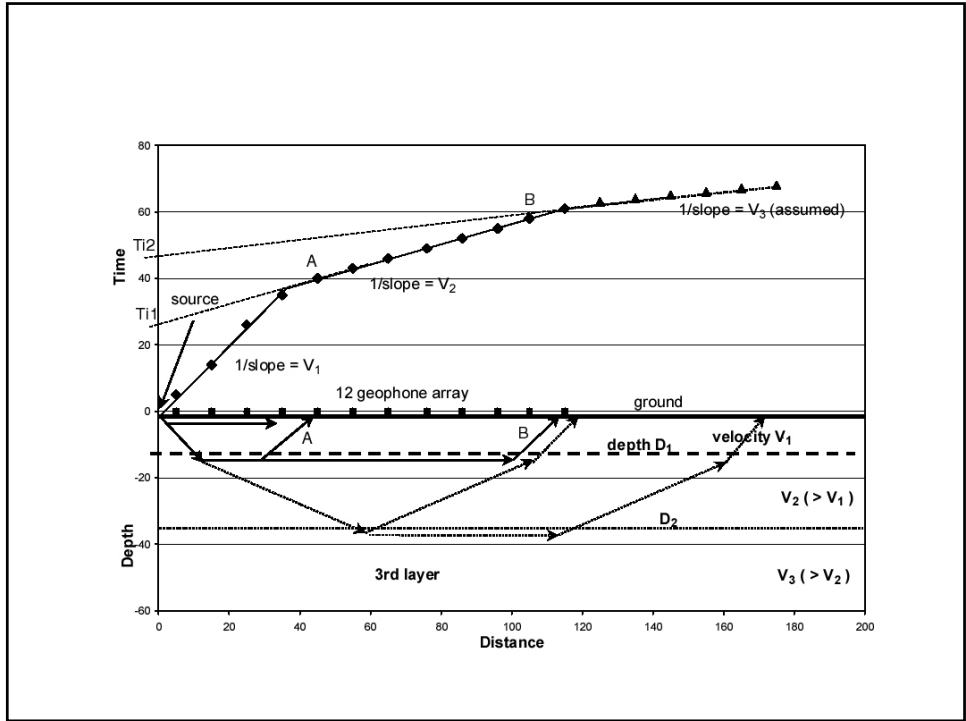


Seismic Refraction

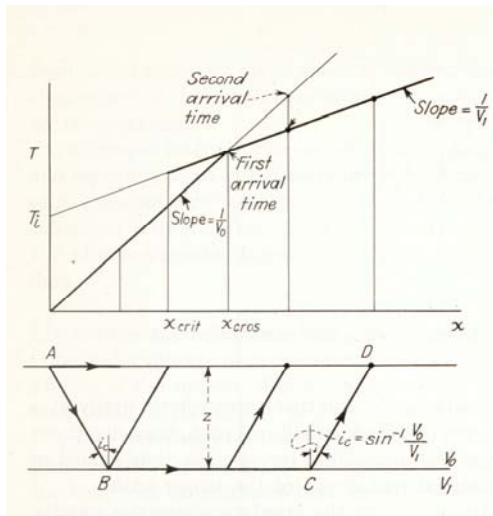








## Two-Layer Refraction Problem

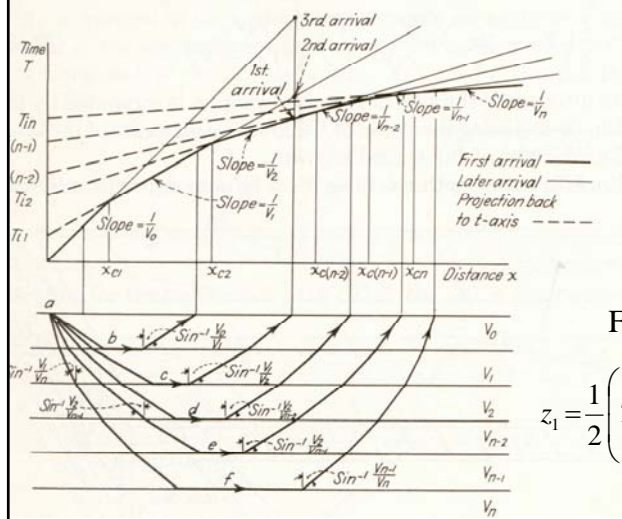


$$X_{crit} = 2z \frac{v_1}{v_0} = 2z \frac{v_1}{\sqrt{v_2^2 - v_1^2}}$$

$$X_C = 2z \sqrt{\frac{v_2 + v_1}{v_2 - v_1}}$$

$$z = \frac{T_{i1} v_0 v_1}{2\sqrt{v_1^2 - v_0^2}}$$

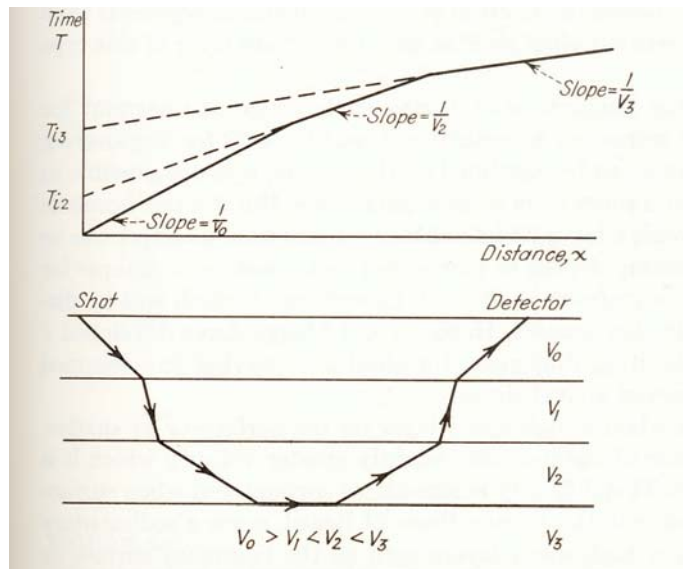
## Multi-layer Problems



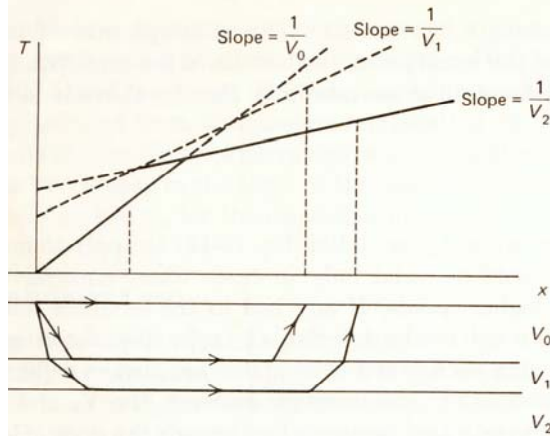
For three layers:

$$z_1 = \frac{1}{2} \left( T_{i2} - 2z_0 \frac{\sqrt{v_2^2 - v_0^2}}{v_2 v_0} \right) \frac{v_2 v_1}{\sqrt{v_2^2 - v_1^2}}$$

## Low Velocity Layers Complications

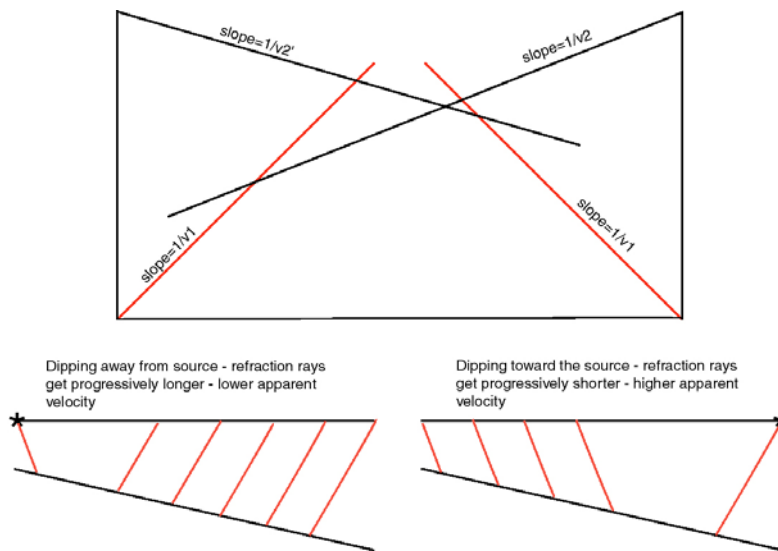


## Hidden Layer Problem

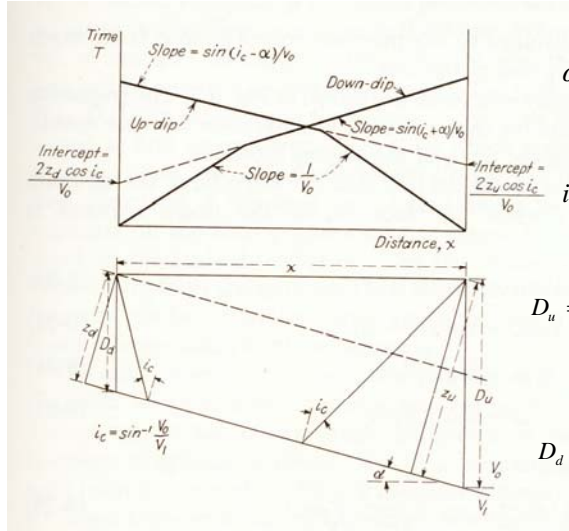


Both low velocity and hidden (thin) layers are complications that cannot be resolved from only first arrival time data. Secondary arrivals or amplitudes need to be employed.

## Dipping Layer Problems



## Dipping Layer Problems



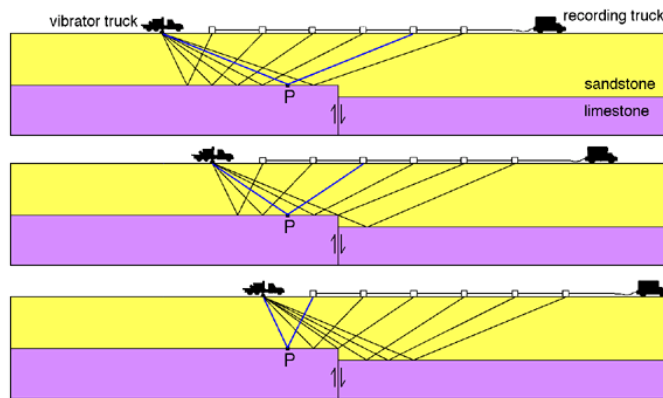
$$\alpha = \frac{1}{2} \left( \sin^{-1} \left( \frac{v_1}{v_d} \right) - \sin^{-1} \left( \frac{v_1}{v_u} \right) \right)$$

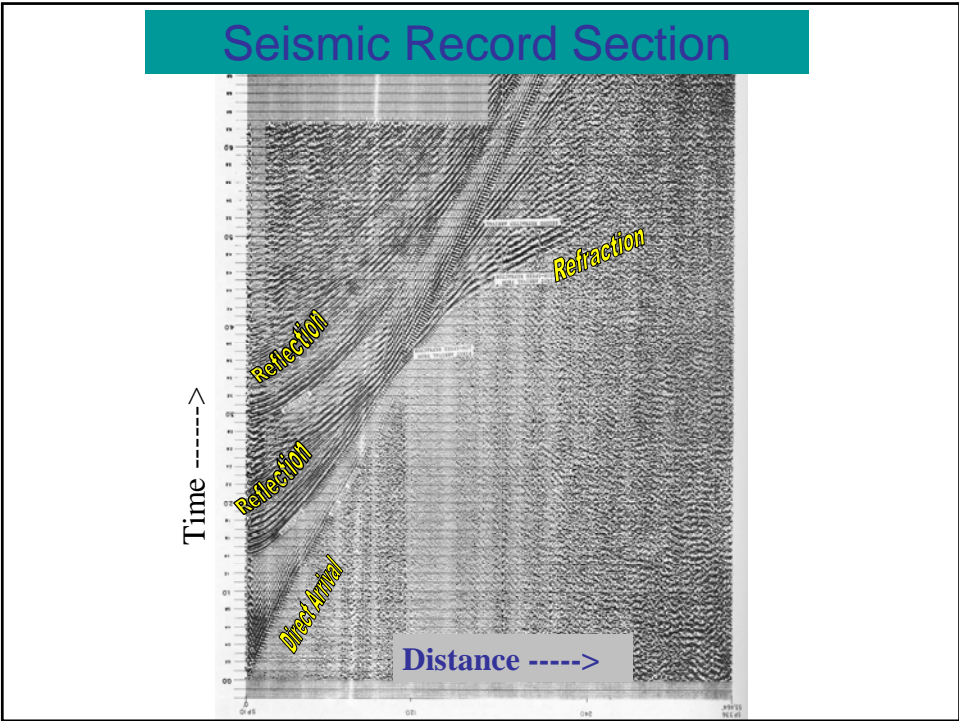
$$i_c = \frac{1}{2} \left( \sin^{-1} \left( \frac{v_1}{v_d} \right) + \sin^{-1} \left( \frac{v_1}{v_u} \right) \right)$$

$$D_u = \frac{T_{iu} v_0 v_1}{2 \cos(\alpha) \sqrt{v_1^2 - v_0^2}} = \frac{T_{iu} v_0}{2 \cos(\alpha) \cos(i_c)}$$

$$D_d = \frac{T_{id} v_0 v_1}{2 \cos(\alpha) \sqrt{v_1^2 - v_0^2}} = \frac{T_{id} v_0}{2 \cos(\alpha) \cos(i_c)}$$

## Seismic Reflection Method





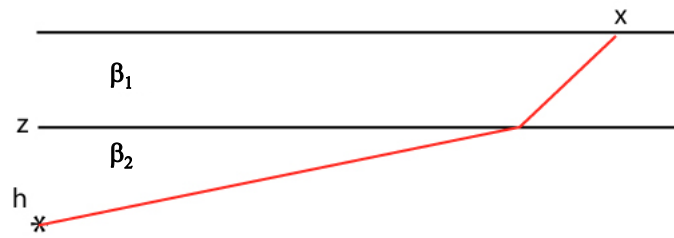
**MODERN  
GLOBAL  
SEISMOLOGY**

THORNE LAY  
TERRY C. WALLACE

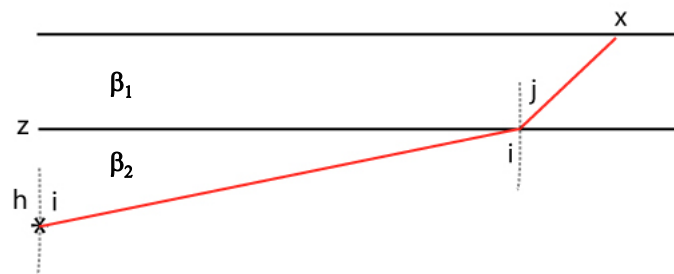
ACADEMIC PRESS

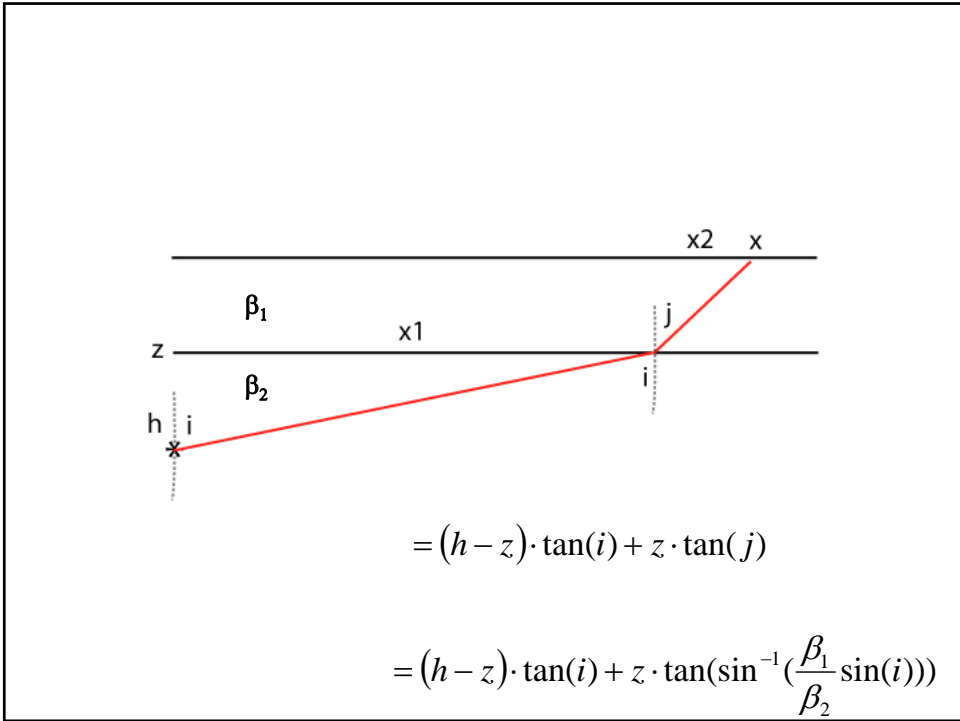
Reflection image courtesy of ChevronTexaco and WesternGeco.

## Homework 4



How do we set up this problem to find arrival time and wave amplitude?





$$x = (h - z) \cdot \tan(i) + z \cdot \tan(\sin^{-1}(\frac{\beta_2}{\beta_1} \sin(i))) \quad \text{Find } i$$

$$p = \frac{\sin(i)}{\alpha_2} \quad \text{Calculate ray parameter}$$

$$\tau = px + \eta_{\alpha_2}(h - z) + \eta_{\alpha_1}z \quad \text{Calculate arrival time}$$

$$T(p) = \frac{2\mu_2\eta_2}{\mu_1\eta_1 + \mu_2\eta_2} \quad \text{Calculate transmission coefficient for amplitude}$$



How about the setup for first multiple arrival?

